

Fig. 3. Lower part: measured slopes  $\partial \ln T_c / \partial p$ . Due to the smallness of the pressure effect on  $T_c$  the error in writing  $\partial \ln T_c / \partial p$  for  $1/T_{c,0} \cdot \partial T_c / \partial p$  is negligible. For p > 25 kbar comparison is made with  $a \cdot \partial \ln \tau / \partial n \equiv a(\partial \ln T_c / \partial n - \partial \ln \theta / \partial n)$  (Eq. (1) in the text). Relevant parameters are plotted in the upper part in the usual nomenclature. Data for  $\theta$  and  $\gamma \propto N$  were taken from Ref.<sup>12</sup> (squares) and <sup>21</sup> (circles). Data points for ternary alloys are marked by triangles

bility)<sup>24</sup>. The value of  $\partial \ln \theta / \partial p$  becomes  $0.65 \cdot 10^{-6} \text{ bar}^{-1}$  which equals our experimental term b within experimental error. Under the reasonable assumption that  $\partial \ln \theta / \partial p$  does not change much with composition, Eq. (2) thus reduces to

$$\partial NV/\partial p = a \cdot \partial NV/\partial n \tag{3}$$

with the solution

$$NV = f(a p+n). \tag{4}$$

24 Bridgman, P. W.: The physics of high pressure, p. 160. London: G. Bell & Sons, Ltd. 1958.

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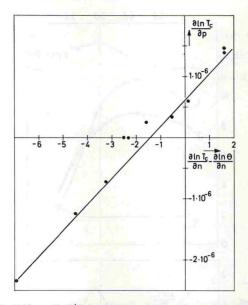


Fig. 4. Plot of  $(\partial \ln T_c/\partial n - \partial \ln \theta/\partial n)$  as obtained from the upper part of Fig. 3, versus  $\partial \ln T_c/\partial p$ . The linear data fit is expressed in Eq. (1)

Thus we arrive at the result that for p > 25 kbar the shape of the function NV(n) remains unaffected by pressure, i.e. is rigid with respect to pressure; within the composition range investigated NV can be changed by the same amount by either adding  $\Delta n$  electrons per atom or applying a pressure  $\Delta p = \Delta n/a$ .

Since in this alloy system V is a slowly varying function of n, the shape of NV(n) is mainly determined by N. One is led to assume that this is also true for NV(p), and to conclude from Eq. (4) that the effect of the application of pressure is primarily a filling of the d-band. A simple mechanism which would provide this can be seen in a pressure-induced shifting of the broad s-band relative to the narrow d-band.

We have used here the original BCS approach. The limited accuracy of the experimental information was thought to inhibit a discussion in terms of more advanced theoretical expressions for  $T_c$  which account for a more realistic model for the Coulomb interaction and the effect of electron-phonon interaction on the electronic properties <sup>9,15,25,26</sup>. Nevertheless, some comments are to be made:

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<sup>25</sup> Morel, P., Anderson, P. W.: Phys. Rev. 125, 1263 (1962).

<sup>26</sup> Garland, Jr., J. W.: Phys. Rev. Letters 11, 111 (1963).